

Biology of the *noisy* gene

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day II: modeling noise

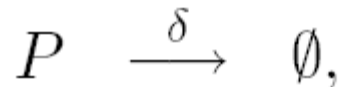
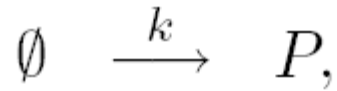


- Two-step gene expression model
- Translational Bursting
- Gillespie's algorithm
- Langevin equations

Stochastic dynamics of gene expression, summary

- Genes are expressed by means of chemical reactions
- Chemical reactions are stochastic processes (collisions, etc)
- Gene expression is noisy: *intrinsic* noise (fluctuating reaction rates) TODAY
extrinsic noise (molecules involved in gene expression)

A simple model of gene expression, summary



Poisson process (birth and death)

- protein produced on average every $1/k$ seconds (**birth**)

$$\frac{d[P]}{dt} = k - \delta[P]$$

- protein decays with rate δ (**death**)

What is the steady state distribution of this model?

What is noise then?

A simple model of gene expression, summary

Recall; a new framework, the master equation

$$\frac{dp(n)}{dt} = -p(n)(r_k + nPr_\delta) + p(n-1)r_k + p(n+1)(n_P + 1)r_\delta$$

Some comments:

- All moments of the distribution $p(n)$ can be derived from it
- It is a linear equation in $p(n)$.
- Solving the master equation can be done for simple systems, however only normally at steady state.
- In connection with experiments, $p(n)$ would represent the fraction of cells having n copies of some given protein

A simple model of gene expression, summary

Poisson distribution

mean $\langle n \rangle = \langle n \rangle_{ss}$

Macroscopic statistics

variance $\sigma^2 = \langle n \rangle_{ss}$

What is noise then?

standard deviation

definition-1 (coe. variat.) =

$$n_1 = \frac{\sigma}{\langle n \rangle}$$

(= $1/\sqrt{\langle n \rangle}$. Poisson distribution, noise increases as the number of molecules decreases)

definition-2 (Fano factor) =

$$n_2 = \frac{\sigma^2}{\langle n \rangle}$$

(= 1, Poisson distribution mean = variance)

Recovering the deterministic law

Equation of the mean; emergence of deterministic law

$$\frac{d\langle n \rangle}{dt} = \sum_n n \frac{dp_n}{dt}$$

$$= \sum_n n [-p_n(r_k + nr_\delta) + p_{n-1}r_k + p_{n+1}(n+1)r_\delta]$$

$$= -r_k \langle n \rangle - r_\delta \sum_n n^2 p_n + r_k \sum_n p_{n-1} n + r_\delta \sum_n n(n+1) p_{n+1}$$

$$= \underline{r_k - r_\delta \langle n \rangle}.$$

Recovering the deterministic law

Considering that $[P] = \frac{\langle n \rangle}{V}$

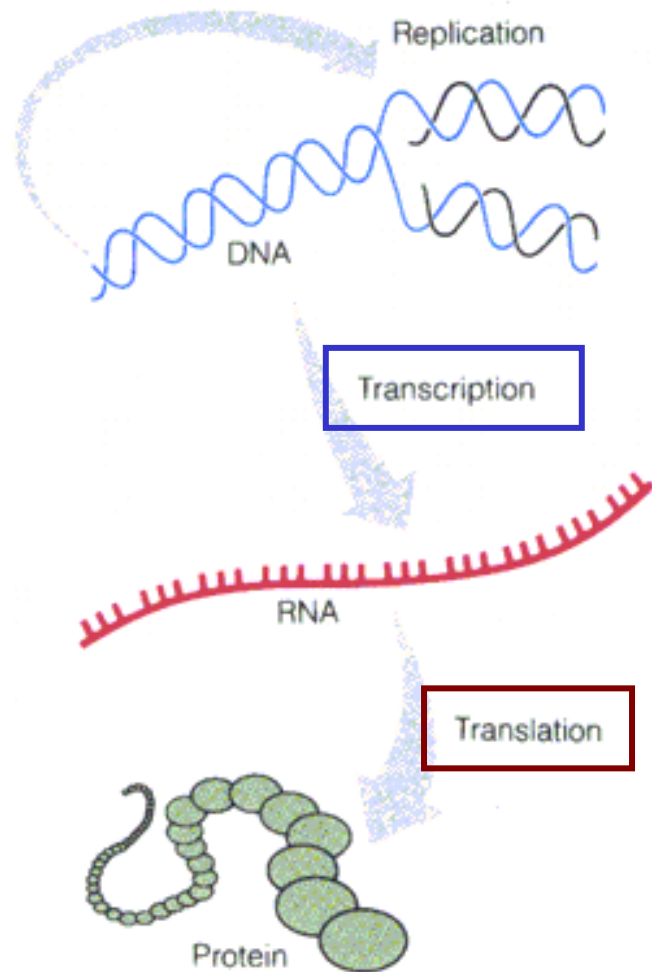
We can rewrite the deterministic equation as

$$\frac{d\langle n \rangle}{dt} = Vk - \delta V [P] = Vk - \delta \langle n \rangle.$$

And thus

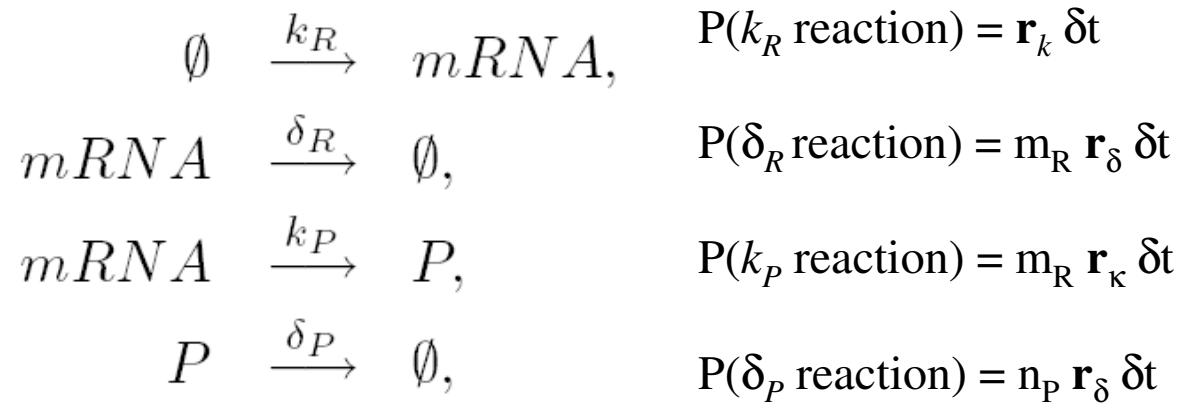
$r_k = Vk$	pseudofirst-order reaction
$r_\delta = \delta$	first order reaction

A more detail model of gene expression



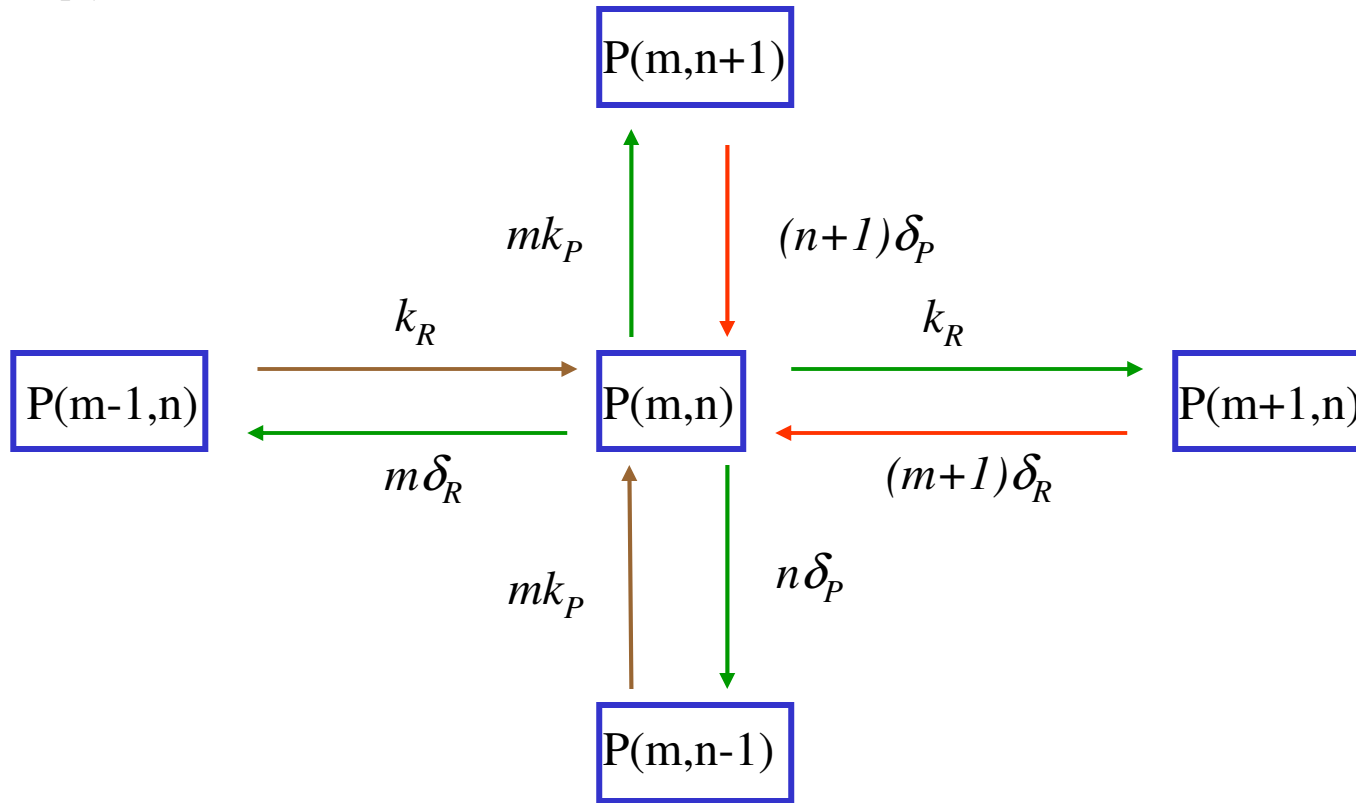
Deterministic model

$$\frac{d[mRNA]}{dt} = k_R - \delta_R[mRNA]$$
$$\frac{d[P]}{dt} = k_P[mRNA] - \delta_P[P]$$



How does the probability of having, say, m mRNA molecules and n P molecules, $p(m,n)$, change with time?

r_k 's as k 's to simplify notation
 (this could also imply that $V = 1$)



$$\begin{aligned} \frac{dp_{m,n}}{dt} &= -p_{m,n}[m\delta_R + mk_P + k_R + n\delta_P] \quad \text{—} \\ &+ p_{m,n+1}(n+1)\delta_P + p_{m+1,n}(m+1)\delta_R \quad \text{—} \\ &+ p_{m,n-1}k_Pm + p_{m-1,n}k_R \quad \text{—} \end{aligned}$$

Equation of the mean; emergence of deterministic laws

note first, a useful equation for a given function $f(n,m)$

$$\begin{aligned} \frac{d\langle f_{n,m} \rangle}{dt} &= -\langle f_{n,m} m \rangle \delta_R - \langle f_{n,m} m \rangle k_P - \langle f_{n,m} \rangle k_R - \langle f_{n,m} n \rangle \delta_P \\ &+ \langle f_{n-1,m} n \rangle \delta_P + \langle f_{n,m-1} m \rangle \delta_R + \langle f_{n+1,m} m \rangle k_P + \langle f_{n,m+1} \rangle k_R \end{aligned}$$

thus, we get

$$\frac{d\langle m \rangle}{dt} = k_R - \delta_R \langle m \rangle \longrightarrow$$

this is the equation the very same equation we obtained for the simple model, i.e., it implies steady state Poisson statistics for **mRNA**

$$\frac{d\langle n \rangle}{dt} = k_P \langle m \rangle - \delta_P \langle n \rangle \longrightarrow$$

what kind of protein macroscopic steady state statistic characterizes **protein** dynamics?

we make use of the following equations ...

$$\frac{d\langle n^2 \rangle}{dt} = -2\langle n^2 \rangle \delta_P + \langle n \rangle \delta_P + 2\langle nm \rangle k_P + \langle m \rangle k_P$$

$$\frac{d\langle nm \rangle}{dt} = -\langle nm \rangle (\delta_P + \delta_R) + \langle m^2 \rangle k_P + \langle n \rangle k_R$$

... to get the final expressions for the macroscopic statistics

$$\text{Fano Protein} = \frac{\langle n^2 \rangle - \langle n \rangle^2}{\langle n \rangle} = 1 + \frac{k_P / \delta_R}{1 + \delta_P / \delta_R} \approx \boxed{1 + \frac{k_P}{\delta_R}} \quad \text{translation efficiency influences noise}$$

$$\text{Fano mRNA} = \boxed{1}$$

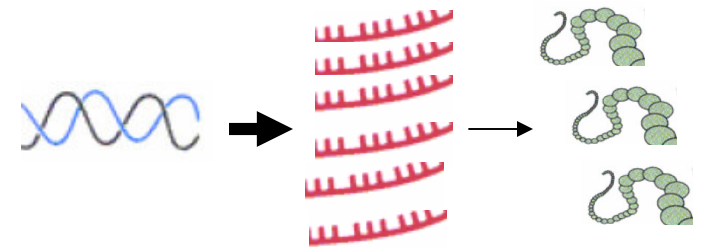
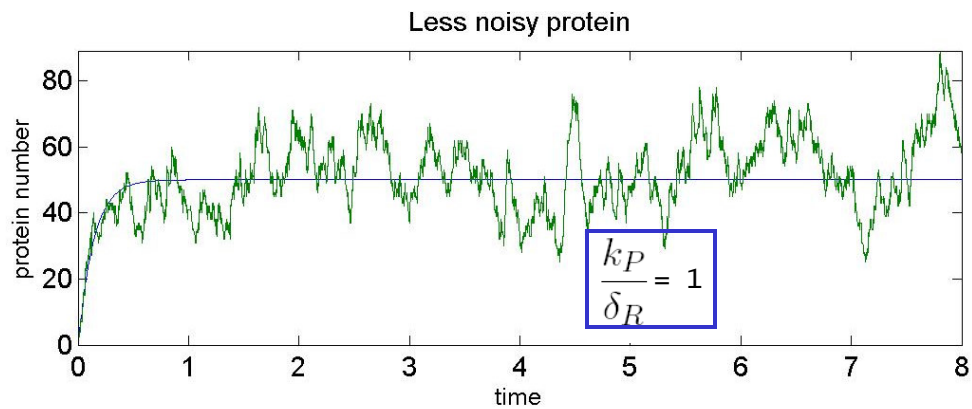
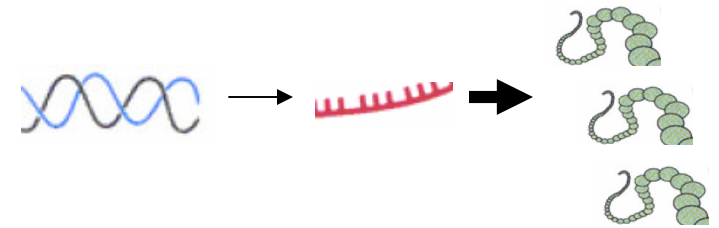
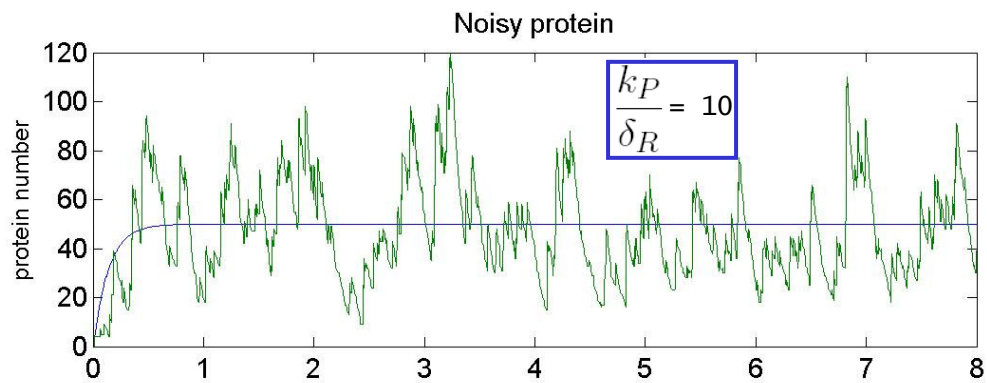


transcription efficiency does not influence noise

protein half-lifetime ~ hours
 mRNA half-lifetime ~ minutes
 thus

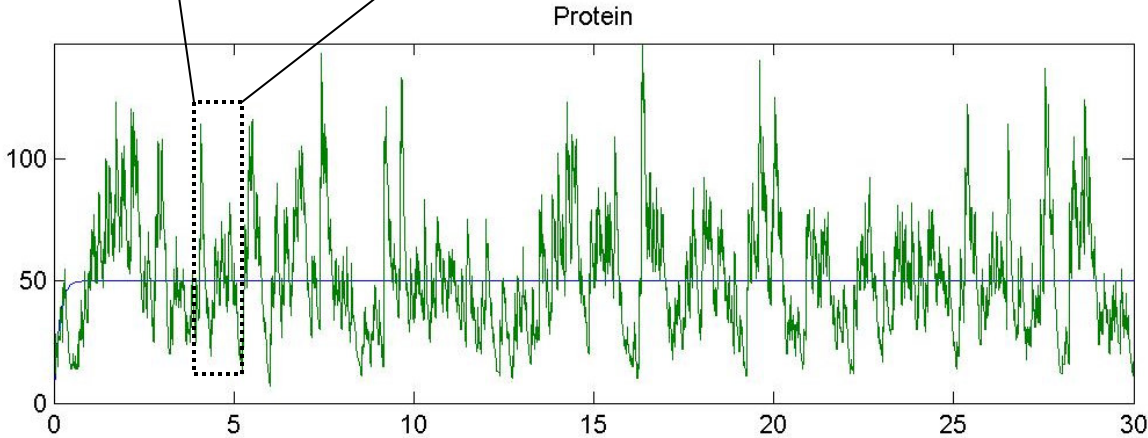
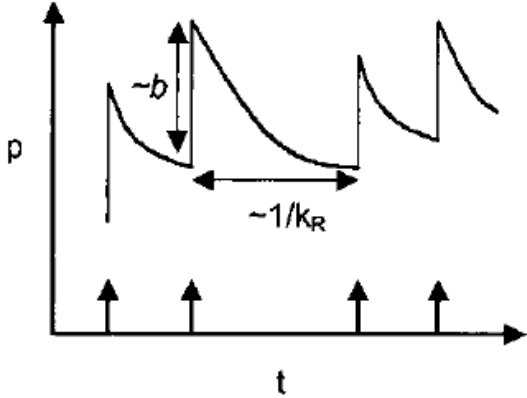
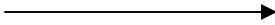
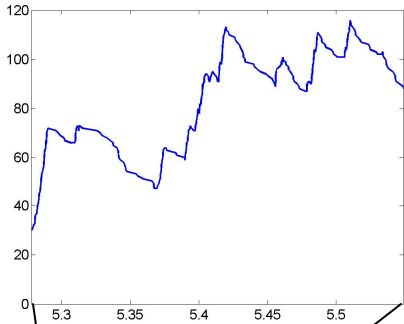
$$t_{1/2} = \log 2 / \delta \quad \text{and} \quad \delta_P \ll \delta_R$$

Independently adjust average and dispersion of protein



“Random bursts model”

$$b = \frac{k_P}{\delta_R}$$



Master equations and gene expression

- Genes are generally regulated by complex nonlinear functions. Analytical studies become difficult.

- Two types of approximation methods

- 1) Numerical Simulation → Gillespie's algorithm

- 2) Perturbation Methods → Langevin equations, ...

Simulating Stochastic Reactions

Two key questions: When will the next reaction occur?
What kind of reaction will it be?

$P(\tau, \mu)d\tau$ = probability that, given the state (X_1, \dots, X_N) at time t , the next reaction in V occurs in the infinitesimal time interval $(t + \tau, t + \tau + d\tau)$ and it will be a R_μ reaction.

propensity function, e.g., $n_1 n_2 \mathbf{r}$

$$P(\tau, \mu)d\tau = P_0(\tau) \downarrow a_\mu d\tau,$$

here,

$P_0(\tau)$: the probability that no reaction happens
in the time interval $(t, t + \tau)$

$a_\mu d\tau$: the probability that reaction
 R_μ will happen in the time interval $(t + \tau, t + \tau + d\tau)$

The function $P_0(t)$: (no reaction)

$$P_0(t+dt) = P_0(t)(1 - a_0 dt) \quad a_0 = \sum_{j=1, M} a_j$$
$$(P_0(t+dt) - P_0(t)) / dt = -a_0 P_0(t)$$

$$d P_0 / dt = -a_0 P_0(t)$$

$$P_0(t) = \exp(-a_0 t)$$

The reaction probability density function:

$$P(\tau, \mu) d\tau = P_0(\tau) a_\mu d\tau = a_\mu \exp(-a_0 \tau) d\tau \quad \mu = 1, \dots, M \quad \tau \in (0, +\infty)$$

It is possible to write $P(\tau, \mu)$ as a product of $P(\tau)$ and $P(\mu)$:

$$P(\tau, \mu) d\tau = a_\mu \exp(-a_0 \tau) d\tau = (a_\mu / a_0) a_0 \exp(-a_0 \tau) d\tau$$

$$P(\mu) = (a_\mu / a_0)$$

$$P(\tau) = a_0 \exp(-a_0 \tau) d\tau$$

Therefore, we may determine the waiting time for the next reaction by generating two random numbers following distributions $P(\tau)$ and $P(\mu)$.

Note that the algorithm is a rigorous consequence of the Fundamental Hypothesis

Gillespie's algorithm

Step 0

Input the desired values for the stochastic rate constants c_1, \dots, c_M . Set the initial molecular population numbers X_1, \dots, X_N and set the time variable t to 0. Initialize the unit-interval random number generator (note UiRN \leftrightarrow distributions $P(\tau)$ and $P(\mu)$).

Step 1

For the current state X_1, \dots, X_N calculate and store M values of propensity functions $a_1 = h_1 c_1, \dots, a_M = h_M c_M$. Accumulate and store the sum of propensity functions $a_0 = \sum_{j=1, M} a_j$

Step 2.

Generate two random numbers $r_1, r_2 \in (0, 1)$ using UiRN. Calculate $\tau = (1/a_0) \ln(1/r_1)$ and take μ to be that integer for which $(a_1 + a_2 + \dots + a_{\mu-1}) < r_2 a_0 \leq (a_\mu + \dots + a_M)$

Step 3.

Update the state of the system by executing one elementary reaction R_μ and increase time of the simulation t by τ .

$t < T_{\max}$

Finish

MATLAB code 2

```
% .. code1stoch.m
% .. simple gene expression stochastic and deterministic

clear all
k = 25;
delta = 1;

% .. stochastic eqs. Gillespie's algorithm
P = 0;
Pstochastic = P;
tmax = 10;
t = 0;
tspan = t;
```

```
while t < tmax

    % .. a's
    a = [k, delta*P(1)];
    a0 = sum(a);
    % .. determine time of next reaction
    r1 = rand;
    tau = -log(r1)/a0;
    t = t + tau;
    % .. determine nature of next reaction
    r2 = rand;
    acumsum = cumsum(a)/a0;
    chosen_reaction = min(find(r2 <= acumsum));

    if chosen_reaction == 1;
        P(1) = P(1) + 1;
    else
        P(1) = P(1) - 1;
    end

    tspan = [tspan,t];
    Pstochastic = [Pstochastic;P];

end
```

```

% .. deterministic eqs.
P0 = 0;
options = [];
[t P] = ode23(@code1equations,tspan,P0,options,k,delta);

```

```

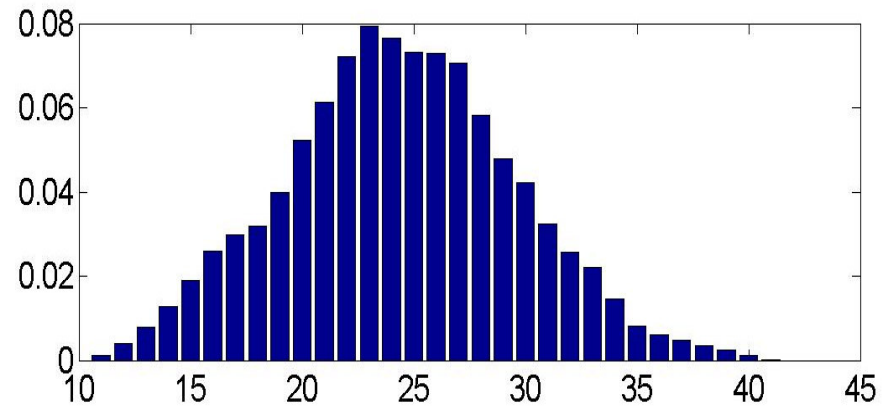
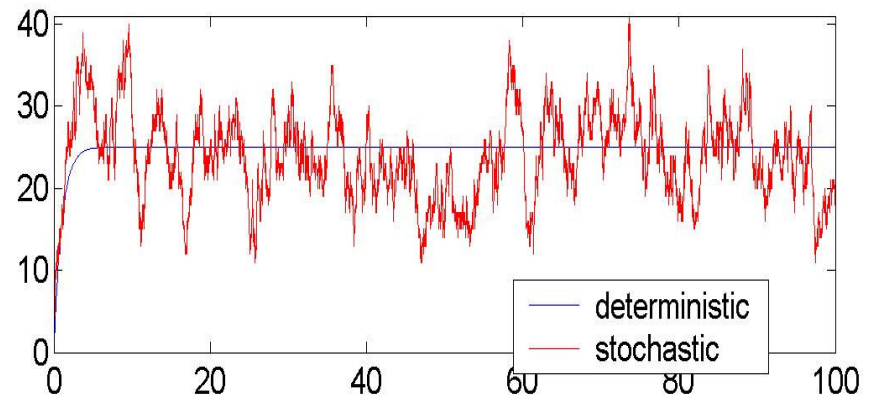
% .. plot
subplot(211)
plot(t,P,t,Pstochastic,'r')
legend('deterministic','stochastic')
axis([0 tmax 0 max(Pstochastic)]);

```

```

subplot(212)
Pst = Pstochastic;
hh = histc(Pst,min(Pst):max(Pst));
bar(min(Pst):max(Pst),hh/sum(hh))

```



meanhist = 24.3909

varihist = 26.7338

fano = 1.0961

Langevin equations

- Developed originally to the study of Brownian motion
- Alternative mathematical framework to that of the Master Equation
- Better suited for an intermediary (“not very noisy”) regime
- Based on adding explicitly noise terms to the deterministic (macroscopic) equations
- Need to characterize the noise distribution of the added noise

$$\frac{d[mRNA]}{dt} = k_R - \delta_R[mRNA] + \underline{\xi_R}$$
$$\frac{d[P]}{dt} = k_P[mRNA] - \delta_P[P] + \underline{\xi_P}$$

ξ_R, ξ_P added stochastic variables.

These equations are fully specified when the probability distributions for the stochastic variables are given.

Valid to describe an intermediate situation where **fluctuations are important even though the number of particles is big enough.**

what are the properties of ξ ?

- we would like to know mean, variance, characteristic fluctuation times, ...

Fluctuation time:

we can ask how much correlates the variation of ξ with respect to its mean:

autocorrelation function

$$\begin{aligned} C_{\xi}(t_1, t_2) &= \left\langle \left[\xi(t_1) - \langle \xi(t_1) \rangle \right] \left[\xi(t_2) - \langle \xi(t_2) \rangle \right] \right\rangle \\ &= \langle \xi(t_1) \xi(t_2) \rangle - \langle \xi(t_1) \rangle \langle \xi(t_2) \rangle \end{aligned}$$

at equal times ($t_1=t_2$) we recover the variance.

Often there exists a characteristic time τ_c for which $C_{\xi}(t_1, t_2) = 0$.

τ_c is known as the **autocorrelation time**

White noise

Langevin originally applied to brownian motion: no reason why thermal fluctuations should favour a particular reaction:

$$\xi_R \quad \xi_P \quad \text{defined such that} \quad \langle \xi(t) \rangle = 0$$

$$\text{thus,} \quad C_\xi(t_1, t_2) = \langle \xi(t_1) \xi(t_2) \rangle$$

collision time is **faster** than time-scale of change of molecule numbers,
noise is uncorrelated

$$C_\xi(t_1, t_2) = \Gamma e^{\frac{-(t_1 - t_2)}{\tau}}$$

↑
noise strength, variance at equal times

$$\tau \rightarrow 0$$

very small autocorrelation times:

$$C_\xi(t_1, t_2) = \Gamma \delta(t_1 - t_2)$$

White noise:

- noise variable with zero autocorrelation time

- white?

$$S(\omega) = \frac{1}{2\pi} \int_{-\infty}^{\infty} e^{-i\omega\tau} \langle \xi(t) \xi(t + \tau) \rangle d\tau = 1$$

all frequencies contribute equally

Color noise

- noise variable with finite autocorrelation time

Stochastic differential equations are very irregular, modified numerical methods to solve them: one to the simplest [Euler-Maruyana method](#)